

Cosmological constraints on the number of neutrino species

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based on work in collaboration with:
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Setting the scene: hints for sterile neutrinos?

Observations at odds with standard 3-neutrino interpretation of global oscillation data

- ◆ LSND anomaly [Aguilar (2001)]
- ◆ MiniBooNE antineutrino results [Aguilar-Arevalo (2010)]
- ◆ Short-baseline reactor experiments
(Bugey, ROVNO, Krasnoyarsk, ILL, Gösgen)
 - ◆ Recent re-evaluation of reactor fluxes [Mention et al. (2011)]

Setting the scene: hints for sterile neutrinos?

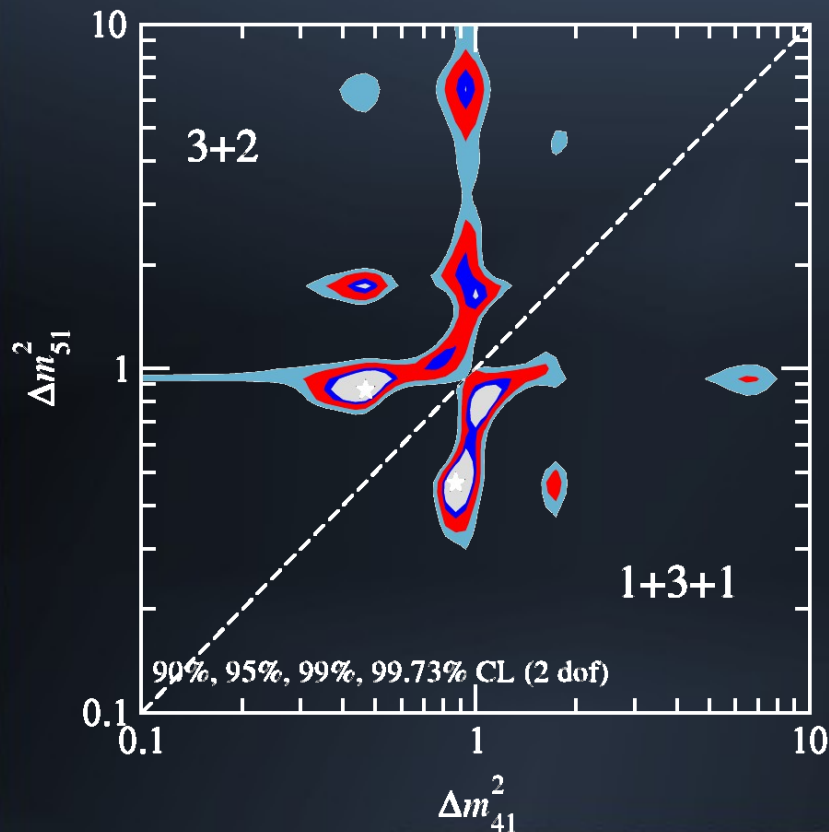
Observations at odds with standard 3-neutrino interpretation of global oscillation data

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Can possibly resolved with oscillations into sterile neutrinos with $\Delta m^2 \sim \text{eV}^2$

Setting the scene: hints for sterile neutrinos?

- Results from reactor/global fit:



| | Δm_{41}^2 [eV ²] | $ U_{e4} $ | Δm_{51}^2 [eV ²] | $ U_{e5} $ | χ^2/dof |
|-----|--------------------------------------|------------|--------------------------------------|------------|---------------------|
| 3+1 | 1.78 | 0.151 | | | 50.1/67 |
| 3+2 | 0.46 | 0.108 | 0.89 | 0.124 | 46.5/65 |

Table I: Best fit points for the 3+1 and 3+2 scenarios from reactor anti-neutrino data.

| | Δm_{41}^2 | $ U_{e4} $ | $ U_{\mu 4} $ | Δm_{51}^2 | $ U_{e5} $ | $ U_{\mu 5} $ | δ/π | χ^2/dof |
|-------|-------------------|------------|---------------|-------------------|------------|---------------|--------------|---------------------|
| 3+2 | 0.47 | 0.128 | 0.165 | 0.87 | 0.138 | 0.148 | 1.64 | 110.1/130 |
| 1+3+1 | 0.47 | 0.129 | 0.154 | 0.87 | 0.142 | 0.163 | 0.35 | 106.1/130 |

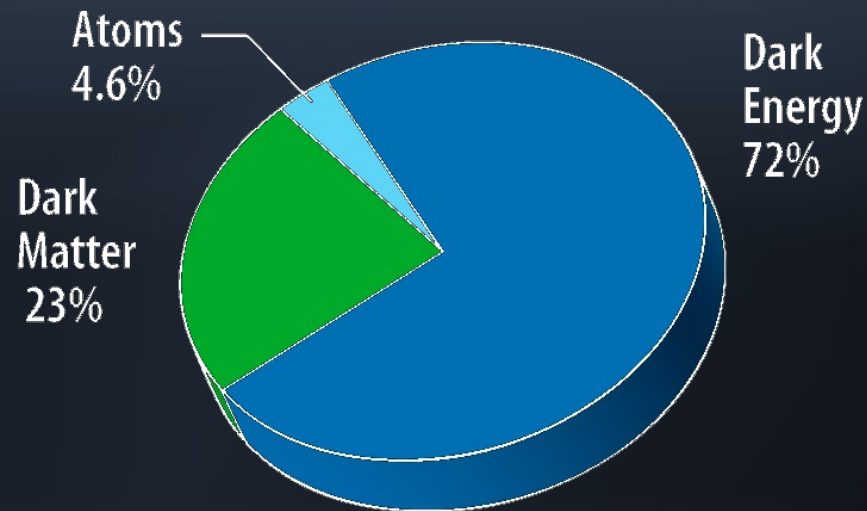
Table II: Parameter values and χ^2 at the global best fit points for 3+2 and 1+3+1 oscillations (Δm^2 's in eV²).

[Kopp, Maltoni, Schwetz (2011)]

What is the Universe made of?

Assuming the Λ CDM-model:

NASA's cosmic pie

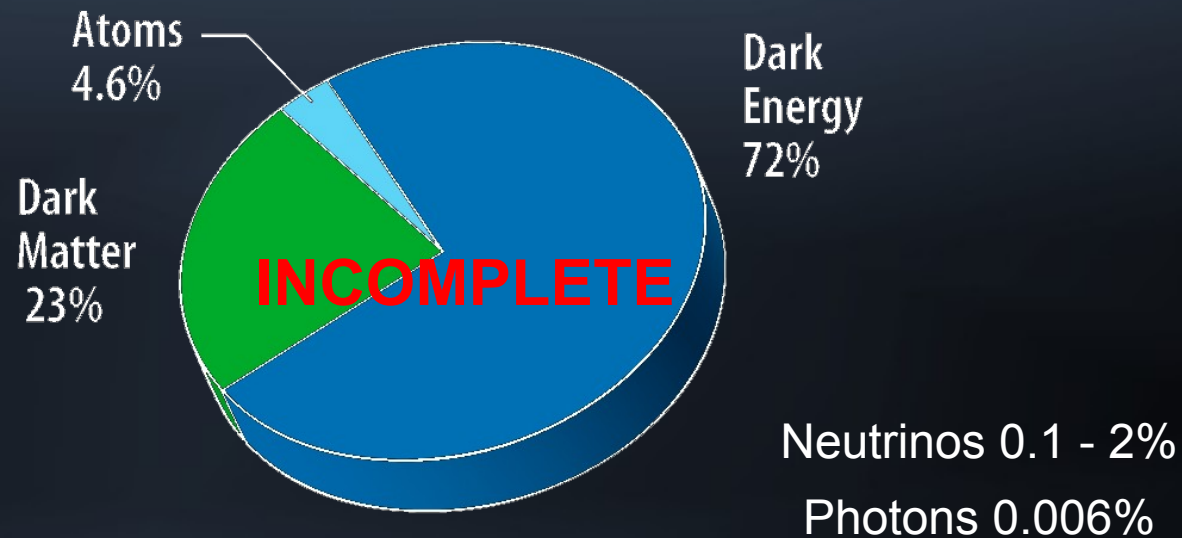


today ($z = 0$)

What is the Universe made of?

Assuming the Λ CDM-model:

NASA's cosmic pie

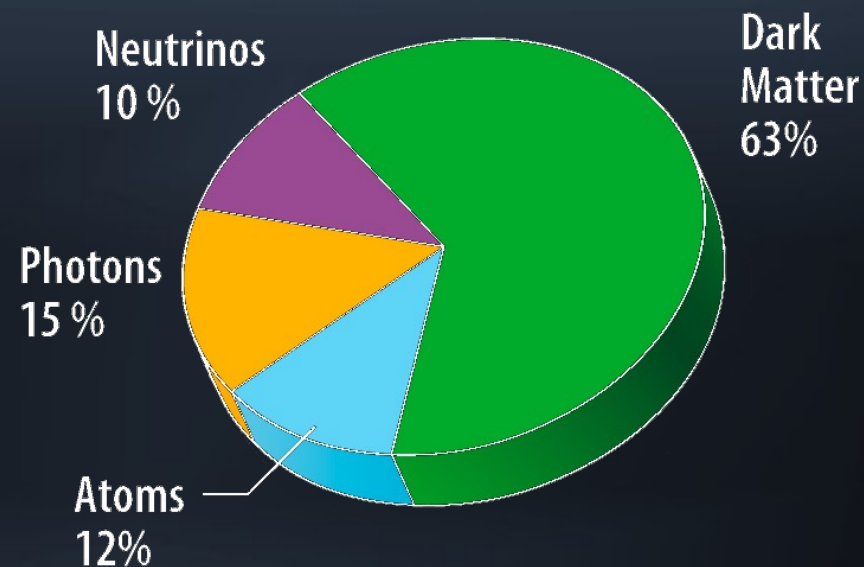


today ($z = 0$)

What is the Universe made of?

Assuming the Λ CDM-model:

NASA's cosmic pie (2)



at decoupling ($z = 1100$)

Radiation content of the Universe

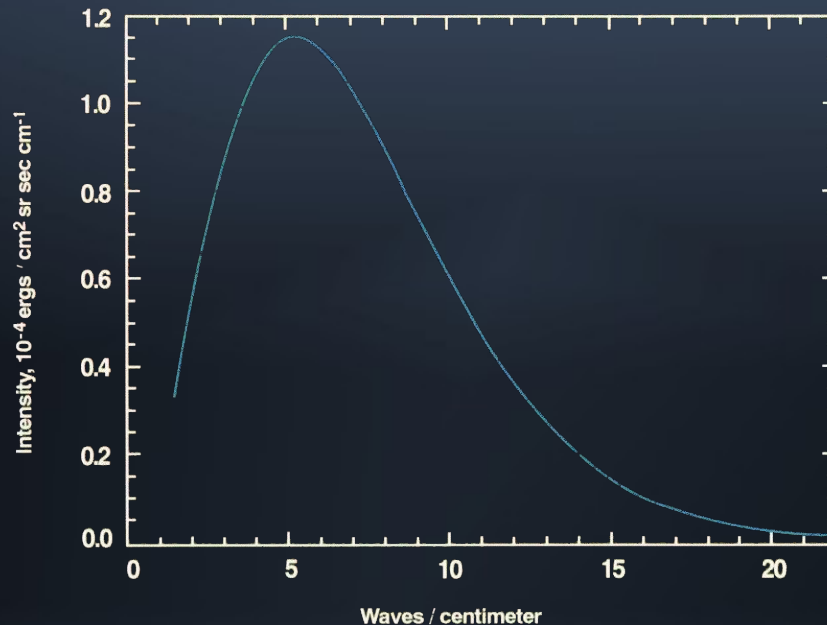
- ♦ Photons: CMB
- ♦ Neutrinos: ν B
- ♦ Other light particle species?

How can we find them?

- ♦ Directly: via scattering
- ♦ Indirectly: via gravitational effects

Cosmic Microwave Background

Directly measured by COBE/FIRAS



[Mather et al. (1993)]

Blackbody spectrum with $T_{\gamma} = 2.725 \pm 0.001 \text{ K}$

[Fixsen & Mather (2002)]

Cosmic Microwave Background

Also affects expansion rate through
photon energy density:

$$\rho_\gamma = \frac{g_\gamma}{(2\pi)^3} \int d^3q \, q \, f_{\text{BE}}(q) = \frac{\pi^2}{15} T_\gamma^4$$

Cosmic Neutrino Background

Neutrinos decouple before e^+e^- -annihilation

$$\longrightarrow T_\nu = \left(\frac{4}{11}\right)^{1/3} T_\gamma \simeq 1.95 \text{ K}$$

- ♦ extremely low energy
- ♦ no direct detection to date

Cosmic Neutrino Background

Neutrino energy density:

$$\rho_{\nu}^{\text{act}} = 3 \cdot \frac{g_{\nu}}{(2\pi)^3} \int d^3q \, q \, f_{\nu}(q) = N_{\text{eff}}^{\text{act}} \cdot \frac{7\pi^2}{120} \left(\frac{4}{11}\right)^{4/3} T_{\gamma}^4$$

LEP: 2.984 ± 0.008

Large mixing ensures that
different mass/flavour eigenstates typically
share a common momentum distribution

[Dolgov et al. (2002), Wong (2002)]

Cosmic Neutrino Background

Neutrino energy density:

$$\rho_{\nu}^{\text{act}} = 3 \cdot \frac{g_{\nu}}{(2\pi)^3} \int d^3q \, q \, f_{\nu}(q) = N_{\text{eff}}^{\text{act}} \cdot \frac{7\pi^2}{120} \left(\frac{4}{11}\right)^{4/3} T_{\gamma}^4$$

For $f_{\nu} = f_{\text{FD}}$, one would have $N_{\text{eff}}^{\text{act}} = 3$

- ♦ Small correction due to ν_e s not being quite completely decoupled at e^+e^- -annihilation + QED correction

————→ Standard Model expectation:

$$N_{\text{eff}}^{\text{act}} = 3.046$$

[Mangano et al. (2005)]

Radiation content of the Universe

Other light stuff?

$$\rho_X = N_X \cdot \frac{7\pi^2}{120} \left(\frac{4}{11}\right)^{4/3} T_\gamma^4$$

Radiation content of the Universe

Other light stuff?

$$\rho_X = N_X \cdot \frac{7\pi^2}{120} \left(\frac{4}{11}\right)^{4/3} T_\gamma^4$$

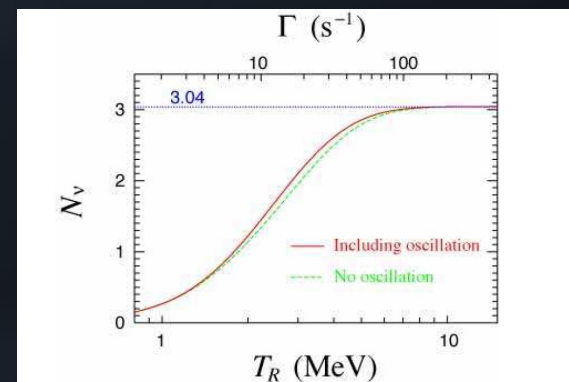
Putting it all together:

$$\begin{aligned} \rho_r &= \rho_\gamma + \rho_\nu^{\text{act}} + \rho_X \\ &= \frac{\pi^2}{15} T_\gamma^4 \left[1 + \underbrace{(N_{\text{eff}}^{\text{act}} + N_X)}_{N_{\text{eff}}} \cdot \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} \right] \end{aligned}$$

A few remarks on N_{eff}

- ♦ is not a constant, in general
 - ♦ increase through light decay products of massive particle
 - ♦ decrease when particles go non-relativistic
 - ♦ (in fact, technically $N_{\text{eff}} \leq 1$ today)

♦ N_{eff} can be < 3.046 at early times if neutrinos out of equilibrium; e.g., low reheating temperature:



[Ichikawa, Kawasaki, Takahashi (2005)]

Determining N_{eff} from observation

Big Bang Nucleosynthesis

- ◆ Primordial element abundances

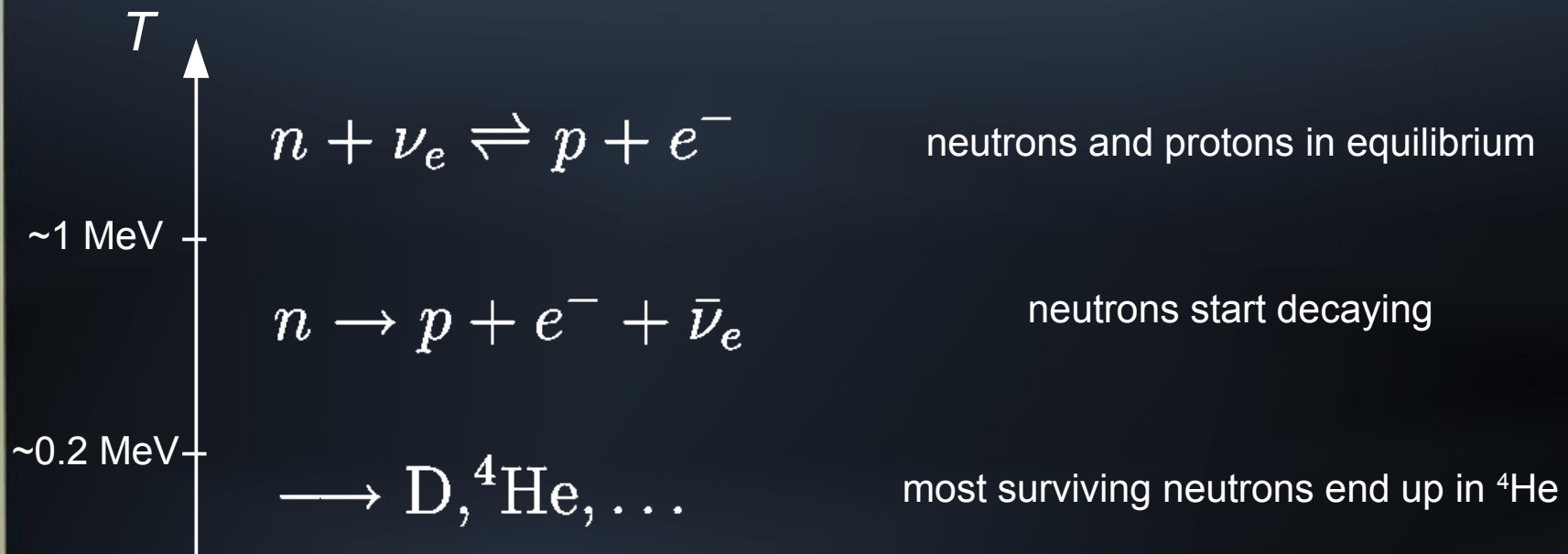
Decoupling

- ◆ Cosmic Microwave Background anisotropies
- ◆ Large scale structure

BBN

Boltzmann equation

nuclear interaction rates \longleftrightarrow expansion rate



BBN

Boltzmann equation

nuclear interaction rate \longleftrightarrow expansion rate

$$\Gamma(\omega_b, f_{\nu_e})$$

$$H \propto \sqrt{\rho_r}$$

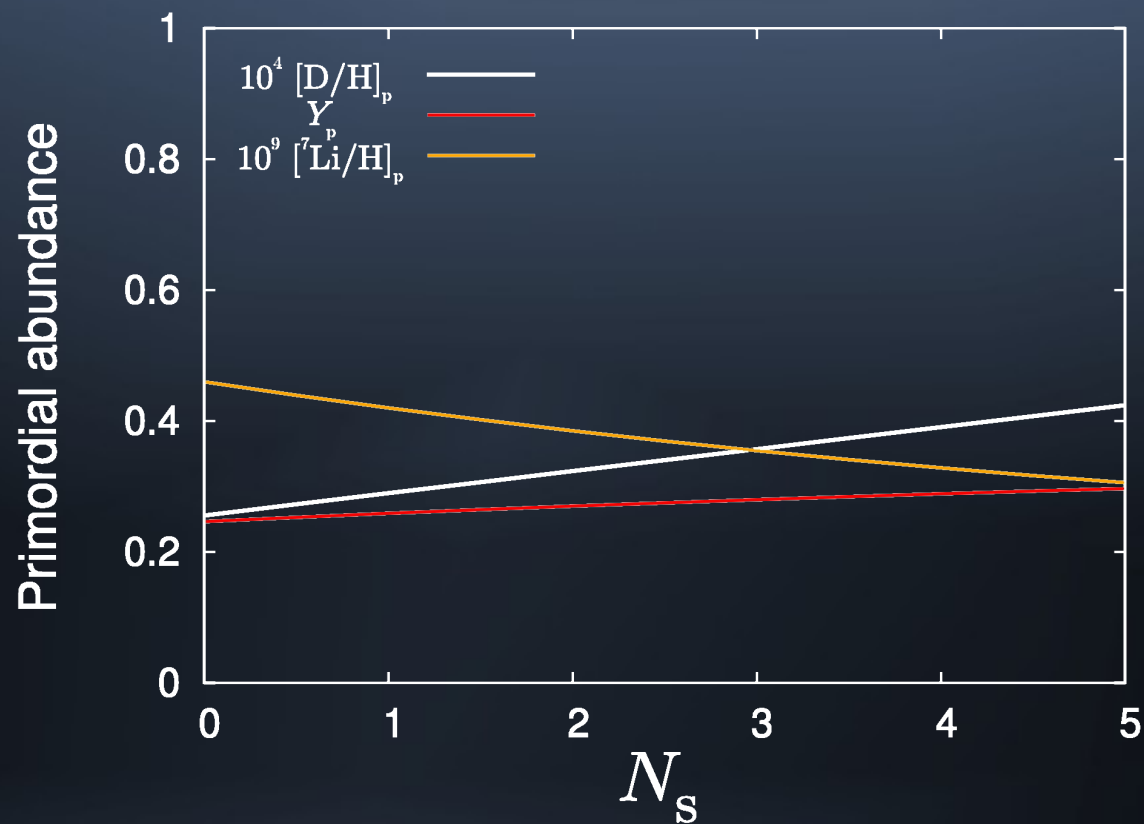
$$N_{\text{eff}}$$

primordial element abundances
as function of $(\omega_b, f_{\nu_e}, N_{\text{eff}}, \dots)$

BBN

- ◆ Assume standard BBN with 3 active neutrinos and N_s additional effective "sterile neutrino" species

BBN



Measure primordial abundances \rightarrow infer N_s

BBN

E.g., Helium:
increasing radiation density
→ higher expansion rate
→ n-p freeze-out at higher T
→ $n/p = \exp[-\Delta m/T]$
→ greater Helium abundance

Measure primordial abundances → infer N_{eff}

Primordial abundances: Deuterium

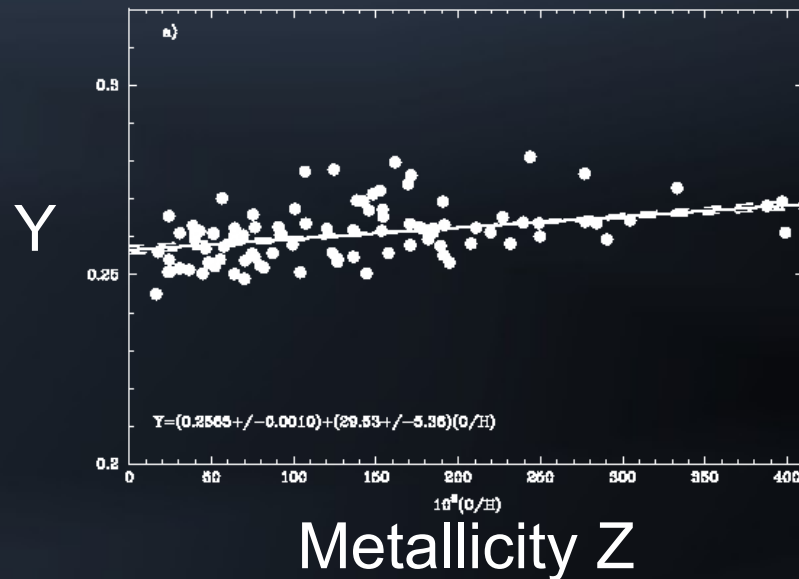
- ◆ Measure absorption of quasar light in low-metallicity hydrogen clouds at high z
- ◆ Relatively "clean" probe
- ◆ Deuterium abundance not subject to strong evolution with redshift
- ◆ From seven measured objects:

$$\log [D/H]_p = -4.55 \pm 0.03$$

[Pettini et al. (2008)]

Primordial abundances: Helium

- ◆ Observe Hydrogen and Helium emission lines in H-II regions of metal-poor dwarf galaxies
- ◆ Astrophysical systematics
 - ◆ Interstellar reddening
 - ◆ Absorption lines in stellar continuum
 - ◆ Radiative transfer
 - ◆ Collisional corrections
- ◆ Helium production by Pop III stars $\rightarrow dY/dZ > 0$

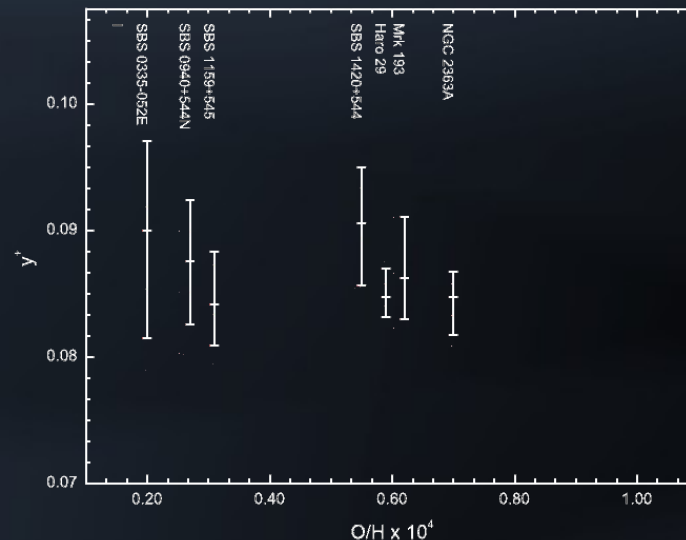


[Izotov, Thuan (2010)]

Primordial abundances: Helium

- ◆ Reliable treatment of systematics for seven "high-quality" systems
- ◆ Linear regression to zero metallicity, limited to positive slopes

$$Y_p = 0.2573^{+0.0033}_{-0.0088}$$



[Aver, Olive, Skillman (2010)]

Calculating theoretical abundances

- ◆ Solve Boltzmann equations numerically (e.g., ParthENoPE)

[Pisanti et al. (2008)]

- ◆ Theoretical uncertainties:

- ◆ Nuclear rates

negligible for Helium, 1.8% for Deuterium

→ folded into likelihood function

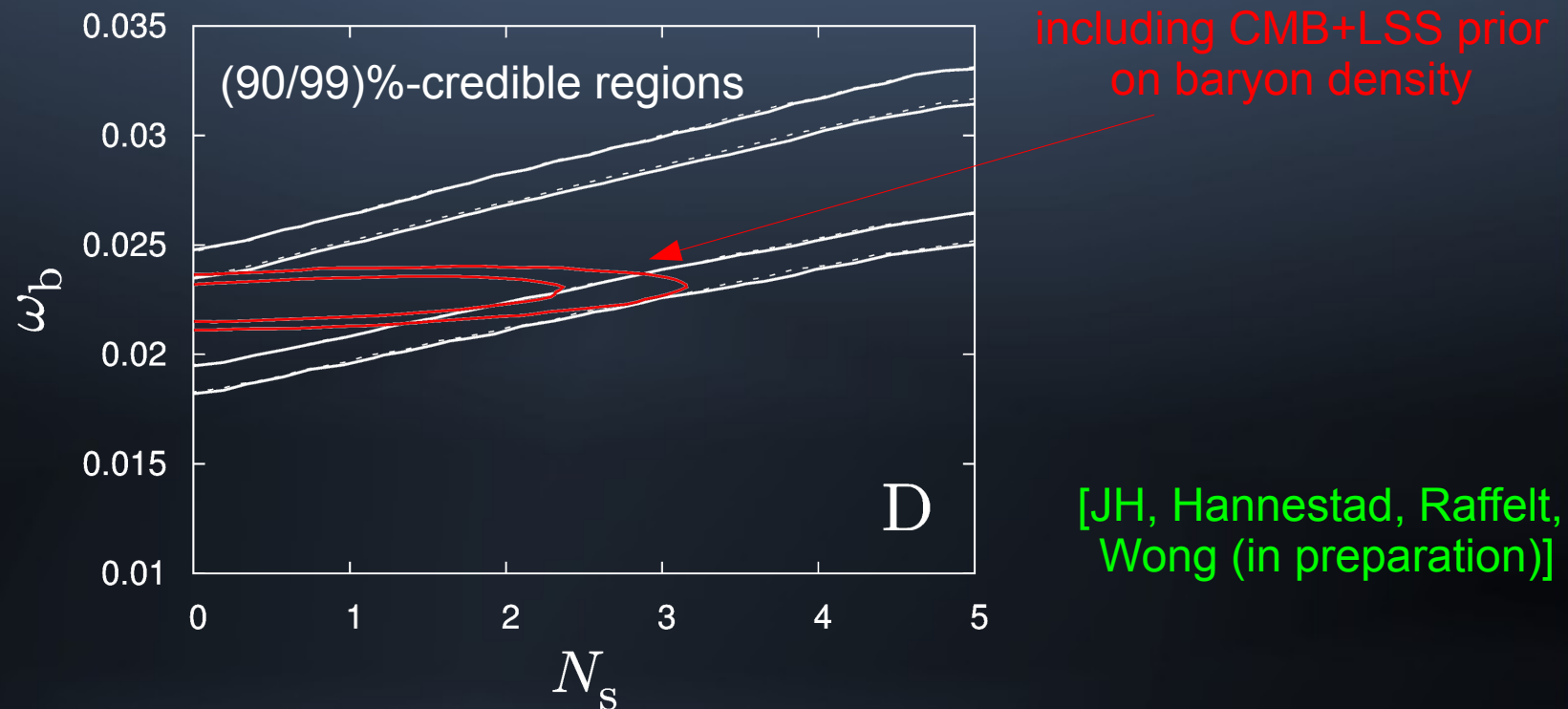
- ◆ Free neutron lifetime

negligible for Deuterium, 0.6% for Helium

Free neutron lifetime

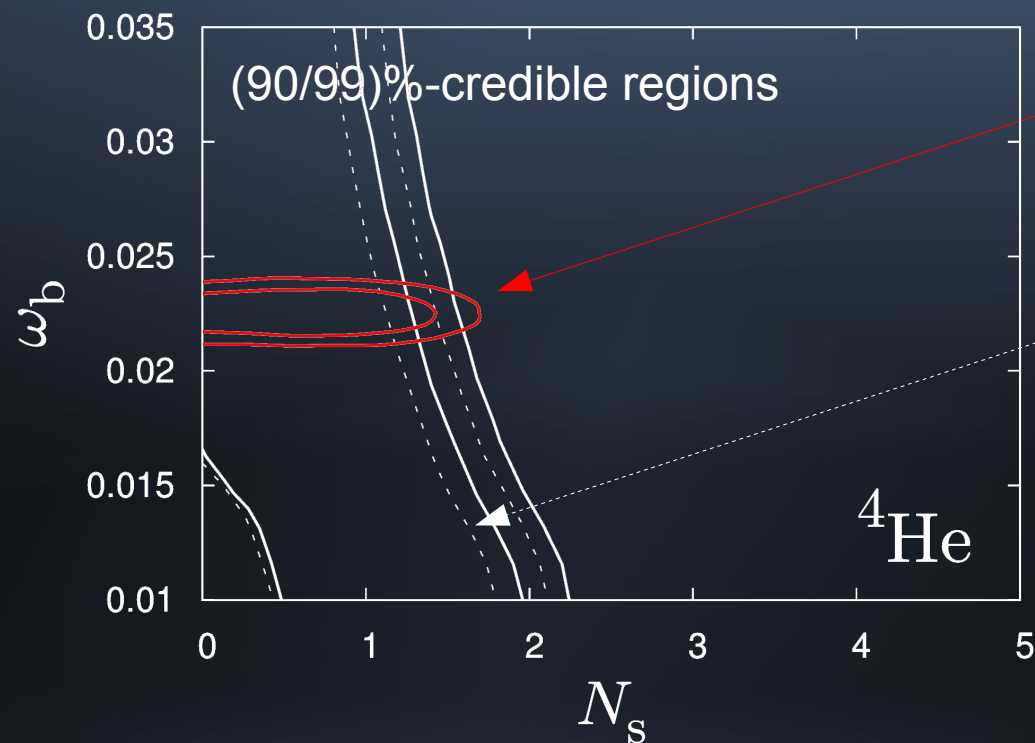
- ◆ No universally agreed upon value:
 - ◆ $\tau_n^{\text{PDG}} = 885.7 \pm 0.8 \text{ s}$ [Particle Data Group (2010)]
 - ◆ $\tau_n^{\text{S}} = 878.5 \pm 0.8 \text{ s}$ [Serebrov et al. (2005)]
 - ◆ $\tau_n^{\text{P}} = 880.7 \pm 1.8 \text{ s}$ [Pichlmaier et al. (2010)]
- ◆ Averaging would not be reasonable
- ◆ Re-analysis of older measurements claims bias of roughly +6 s [Serebrov, Fomin (2010)]
- ◆ We consider the two extremal values until the matter is settled

BBN constraints: Deuterium only



- ◆ N_s unconstrained from Deuterium alone
- ◆ Combination with baryon density prior gives upper limit

BBN constraints: Helium only



including CMB+LSS prior
on baryon density

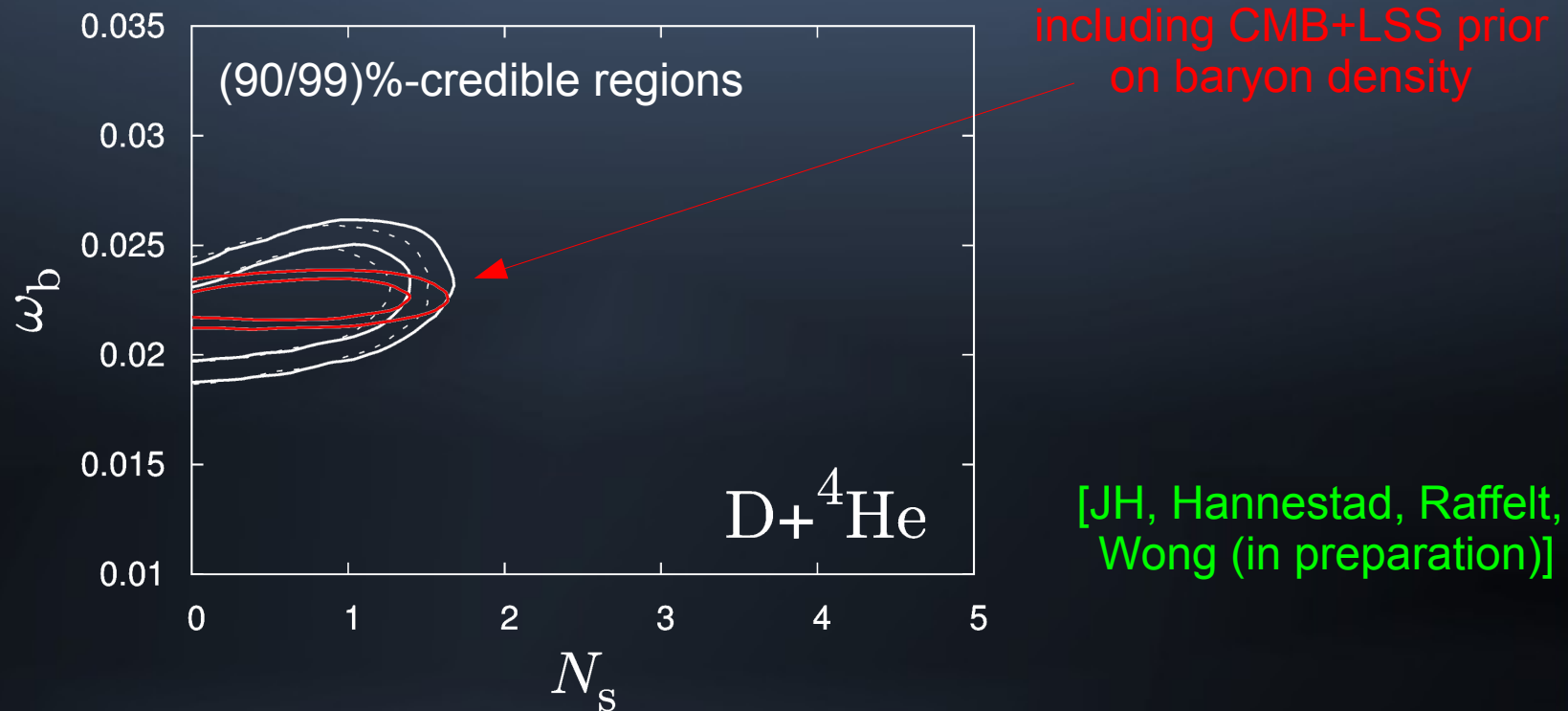
effect of going to higher
neutron lifetime

→ higher predicted Y_p
→ tighter bound on N_s

[JH, Hannestad, Raffelt,
Wong (in preparation)]

- ♦ Helium is a much better probe of N_s

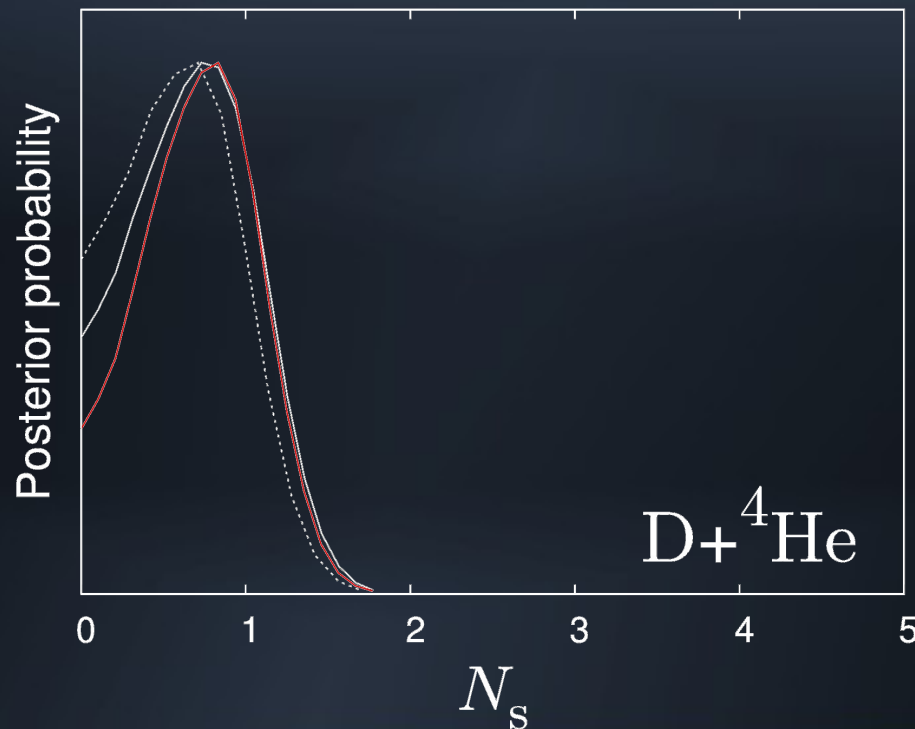
BBN constraints: Deuterium + Helium



- ◆ $N_s < 1.26$ (**1.24**) @95% credibility
- ◆ Best-fit at $N_s = 0.86$

BBN and sterile neutrinos

- ◆ 3+1 scenario slightly preferred over 3+0
- ◆ 3+2 ruled out at high significance ...



[JH, Hannestad, Raffelt,
Wong (in preparation)]

BBN and sterile neutrinos

- ◆ 3+1 scenario slightly preferred over 3+0
- ◆ 3+2 ruled out at high significance ...
... unless:
 - ◆ Incomplete thermalisation
→ effective N_s is smaller than 2
 - ◆ Non-standard BBN?

Degenerate BBN

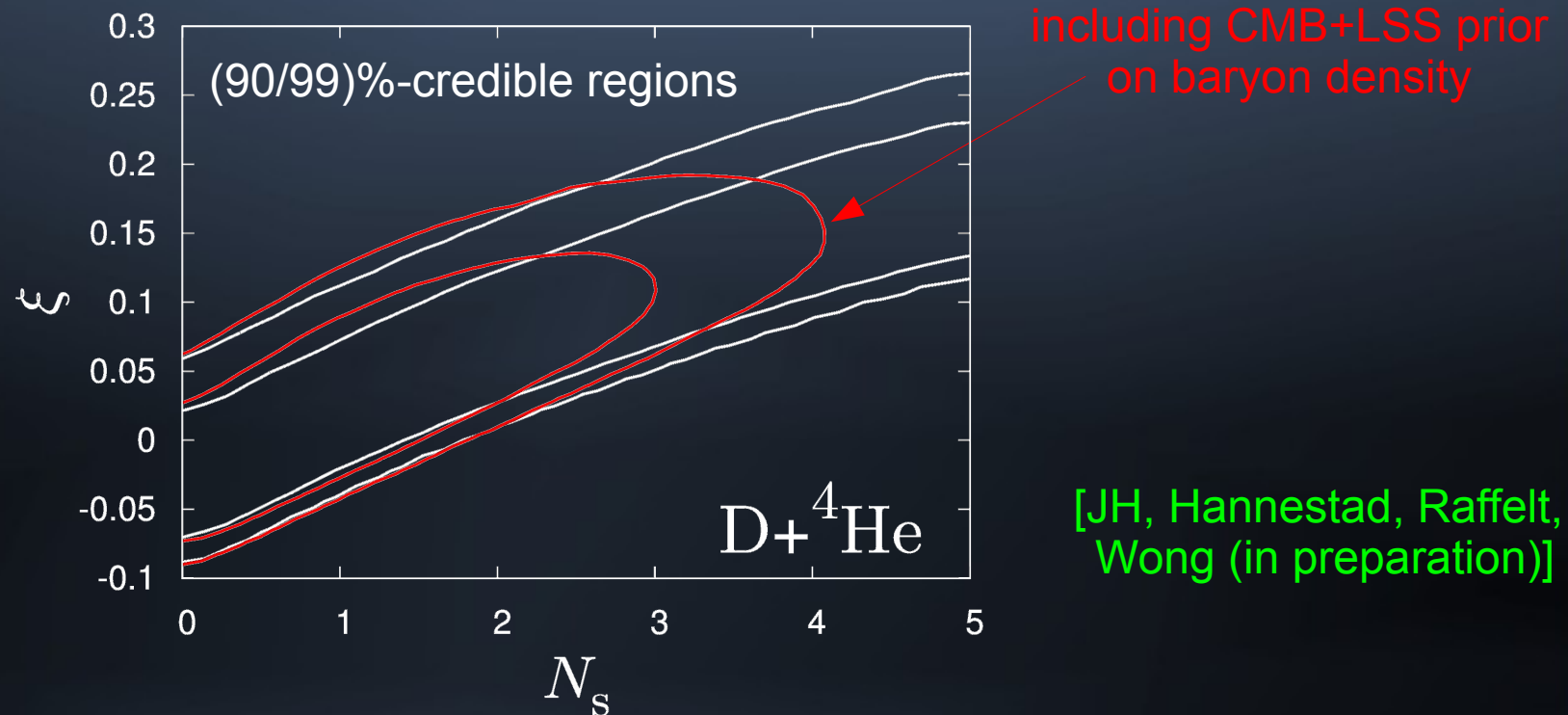
- ◆ Allow for a neutrino chemical potential ξ
- ◆ Assume all active species share the same ξ
- ◆ Two effects:
 - ◆ Additional radiation energy density

$$\Delta N_{\text{eff}} = \frac{45}{7} \left[2 \left(\xi/\pi \right)^2 + \left(\xi/\pi \right)^4 \right]$$

- ◆ Change in initial equilibrium n/p ratio

$$n/p = \exp \left(-\frac{\Delta m}{T} - \xi \right)$$

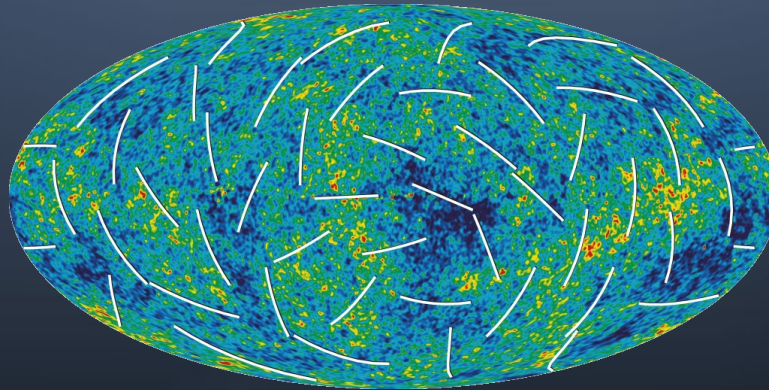
BBN constraints: Degenerate BBN



- ◆ ξ of order 0.1 could save 3+2 (and even 3+3!)

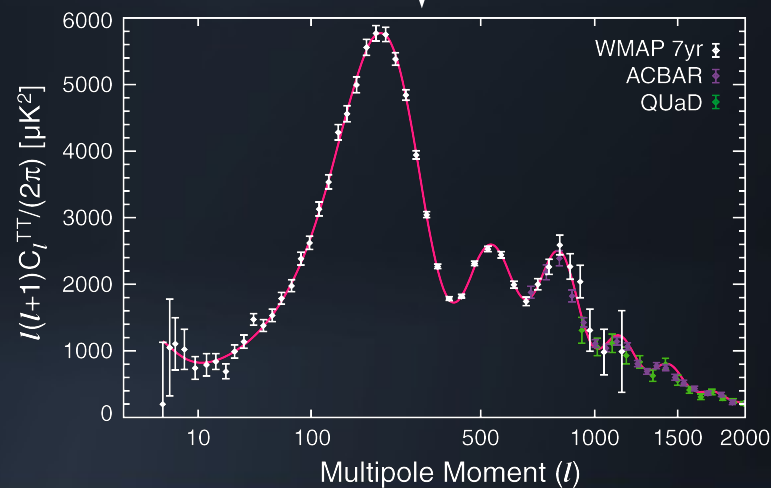
N_{eff} and the CMB

CMB map



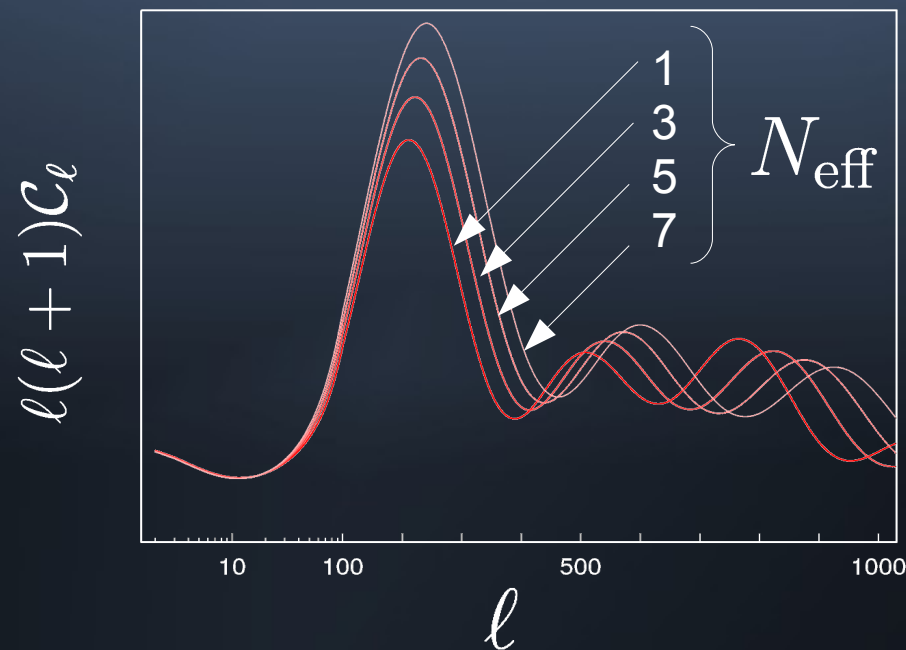
expand in spherical harmonics

CMB angular
power spectrum



[WMAP (2010)]

N_{eff} and the CMB



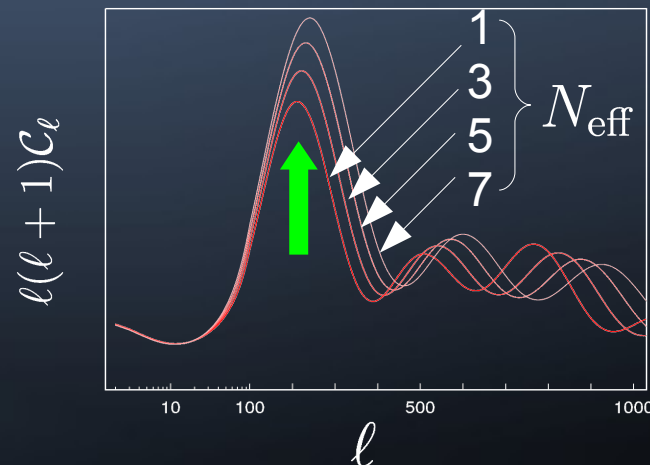
- ♦ Angular power spectrum is a function of $O(10)$ cosmological parameters (e.g., ω_b , ω_{dm} , ω_v , Ω_{de} , N_{eff} , ...)

N_{eff} and the CMB

- ◆ Matter-radiation equality
- ◆ Sound horizon
- ◆ Anisotropic stress
- ◆ Damping tail

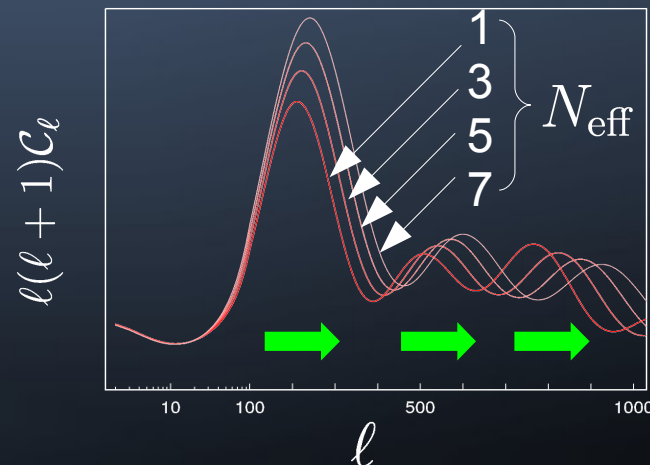
$$1 + z_{\text{eq}} = \frac{\Omega_{\text{m}}}{\Omega_{\text{r}}} \simeq \frac{\Omega_{\text{m}} h^2}{\Omega_{\gamma} h^2} \frac{1}{1 + 0.2271 N_{\text{eff}}}$$

- ◆ Completely degenerate with matter density
- ◆ Larger $N_{\text{eff}} \rightarrow$ later equality \rightarrow enhanced early integrated Sachs-Wolfe-effect \rightarrow higher first peak



N_{eff} and the CMB

- ◆ Matter-radiation equality
- ◆ Sound horizon
- ◆ Anisotropic stress
- ◆ Damping tail

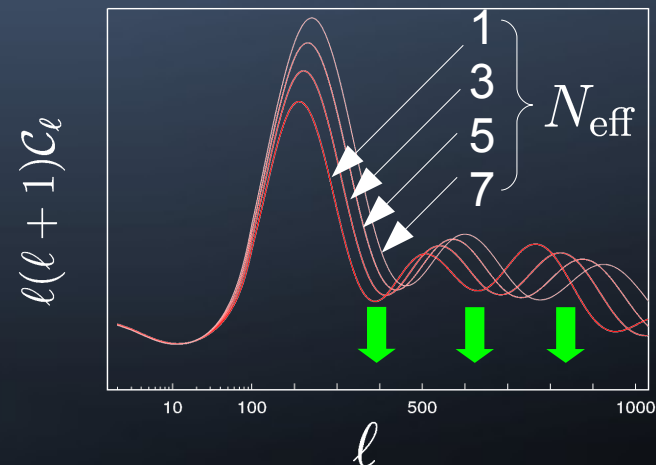


- ◆ Function of radiation, baryon and matter densities
- ◆ θ_s = Sound horizon/distance to last scattering surface
determines positions of acoustic peaks

also depends on dark energy density

N_{eff} and the CMB

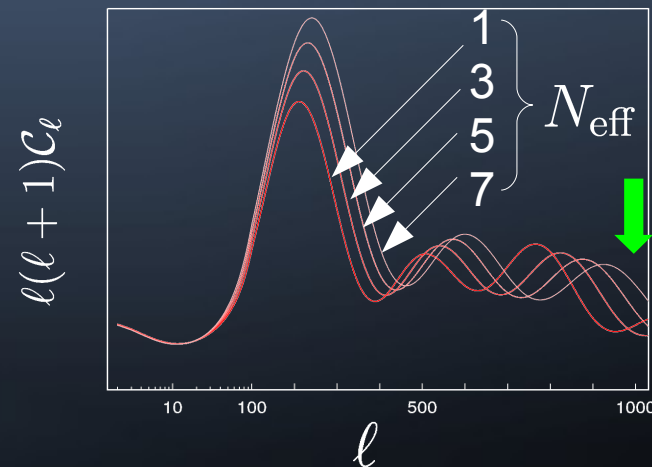
- ◆ Matter-radiation equality
- ◆ Sound horizon
- ◆ Anisotropic stress
- ◆ Damping tail



- ◆ Free streaming particles \rightarrow anisotropic stress
- ◆ Dampens fluctuations during radiation domination
- ◆ Suppression of power at multipoles > 200

N_{eff} and the CMB

- ◆ Matter-radiation equality
- ◆ Sound horizon
- ◆ Anisotropic stress
- ◆ Damping tail



- ◆ Last scattering surface has finite thickness
→ exponential damping of fluctuations below damping scale
- ◆ For fixed peak positions, increasing N_{eff} enhances damping

N_{eff} and the CMB

- ◆ Damping tail can help break degeneracies with other parameters

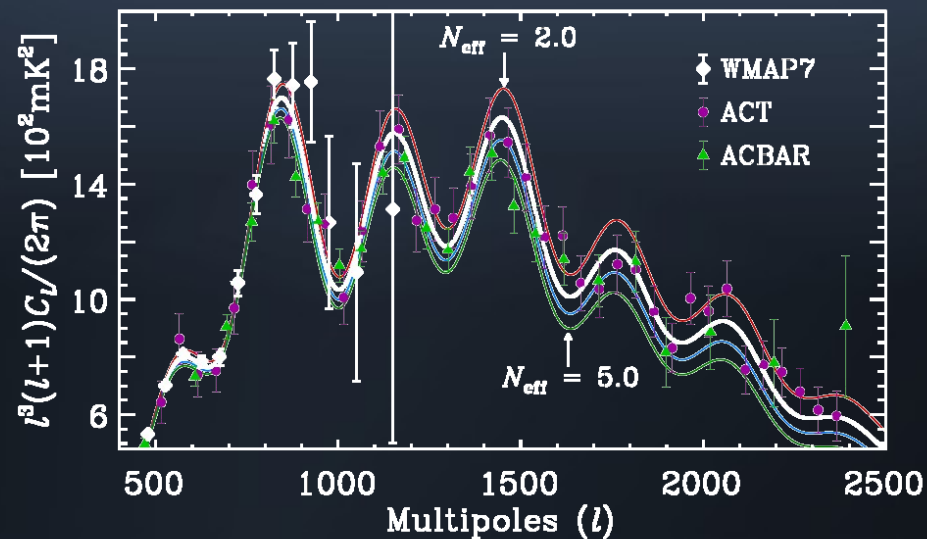
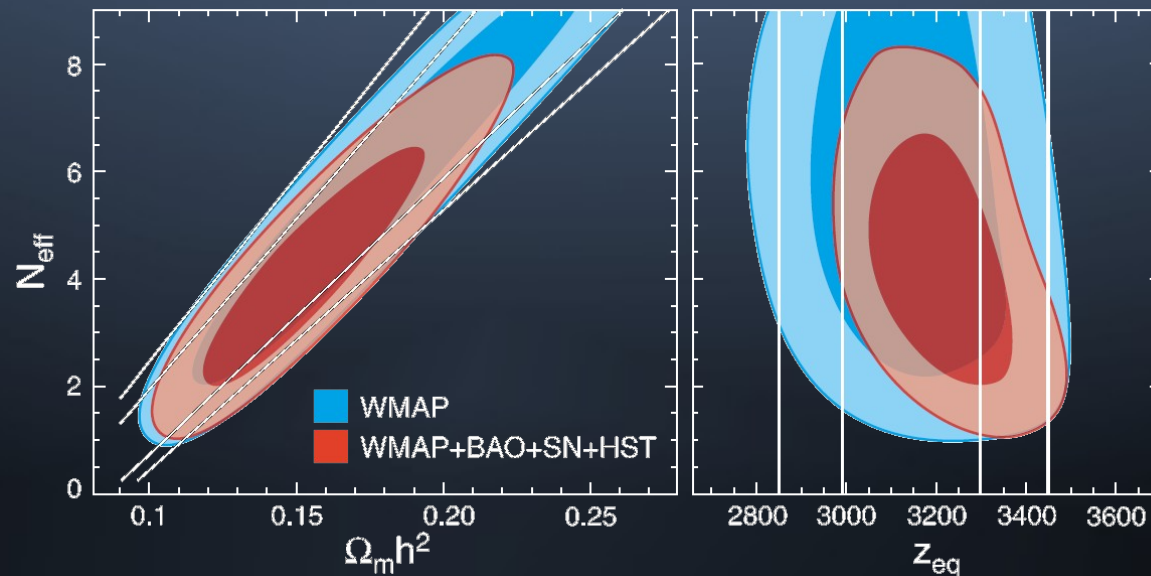


FIG. 1. *Top panel:* WMAP, ACBAR and ACT power spectrum measurements, and theoretical power spectra normalized at $\ell = 200$ for N_{eff} varying from 2 to 5 with ρ_b , θ_s , and z_{EQ} held fixed.

[Millea et al. (2011)]

N_{eff} from WMAP+LSS+...

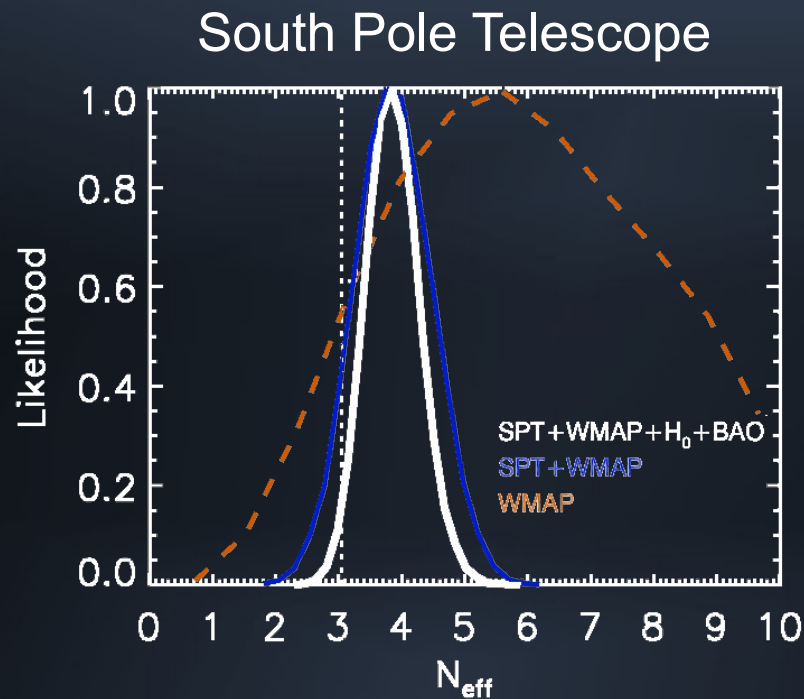


- ♦ lower limit from WMAP alone (\rightarrow anisotropic stress)
- ♦ meaningful upper limit requires combination with other data sets sensitive to matter density and expansion rate ...

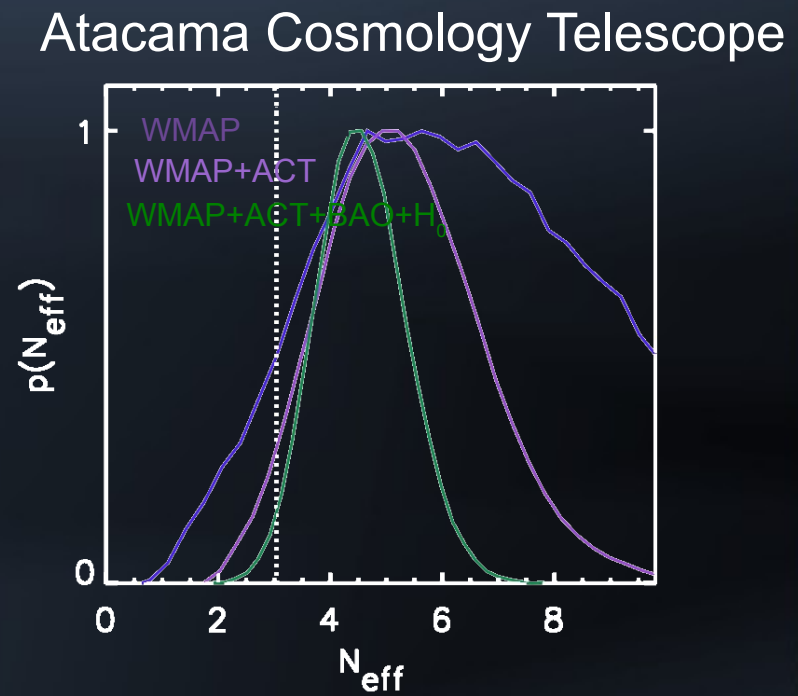
[WMAP: Komatsu et al. (2008)]

N_{eff} from CMB alone

- ... or measurement of the damping tail of the CMB angular power spectrum

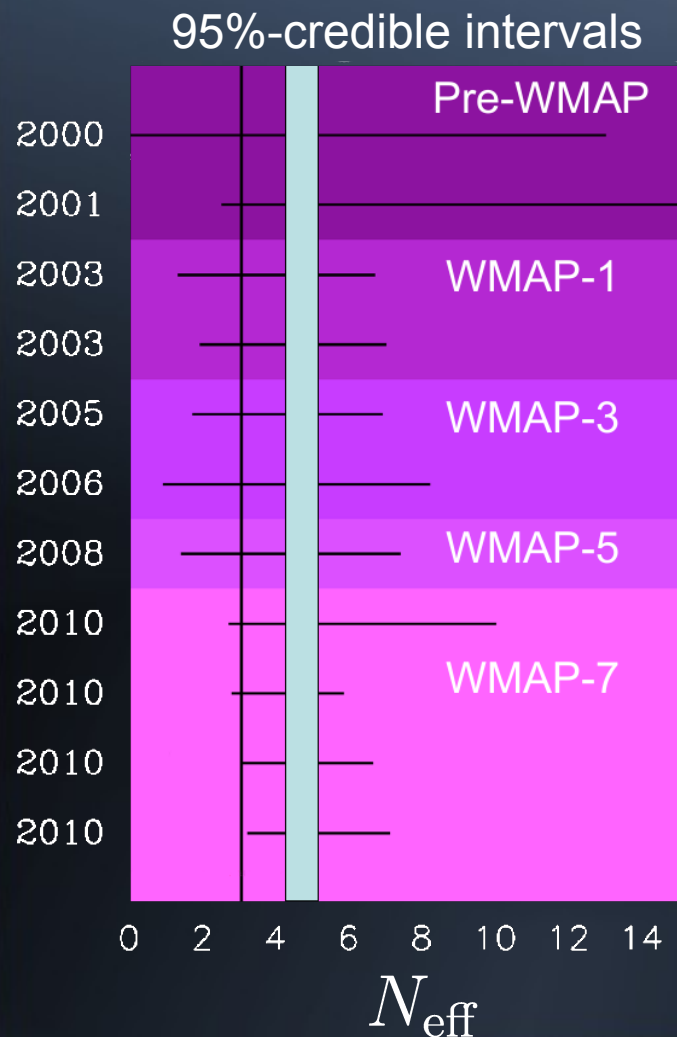


[Keisler et al. (2011)]



[Dunkley et al. (2010)]

CMB+X bounds on N_{eff}



◆ Precise numbers depend on cosmological model and data sets used

◆ Recent analysis: $N_{\text{eff}} = 4.47^{+1.82}_{-1.74}$

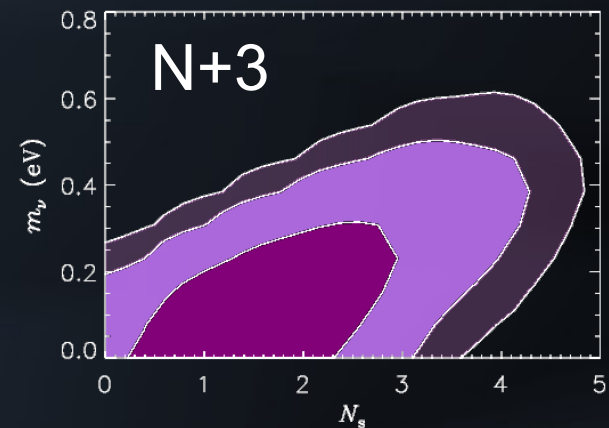
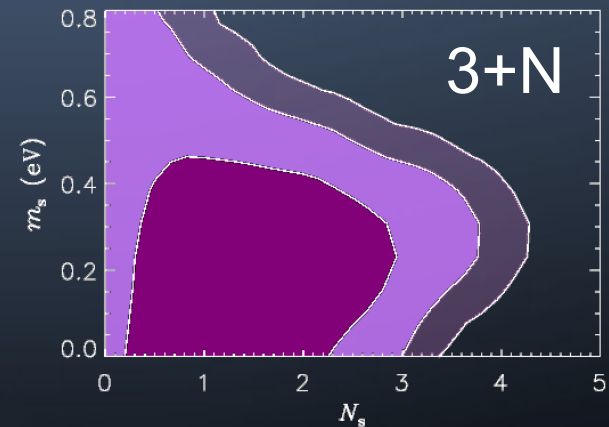
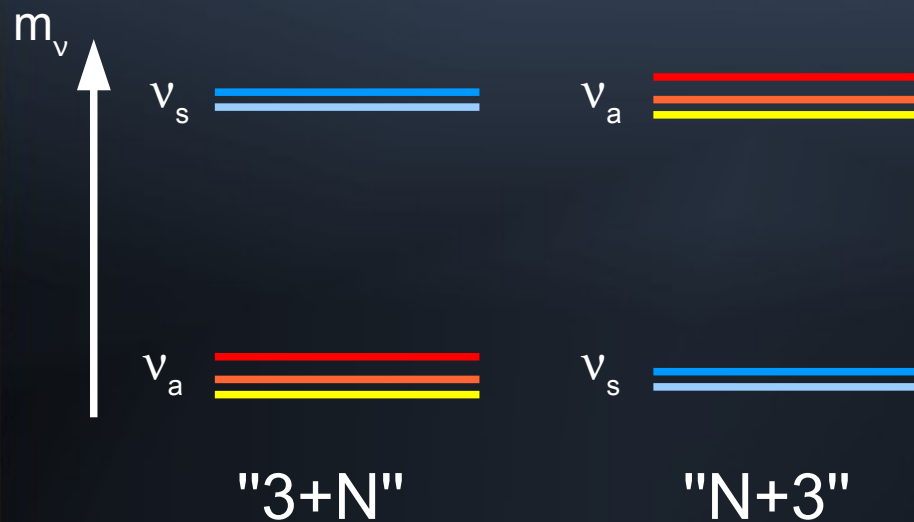
CMB + SDSS-DR7-BAO + HST

Λ CDM + neutrino mass + N_{eff}

[JH, Hannestad, Lesgourgues,
Rampf, Wong (2010)]

Sterile neutrino scenario

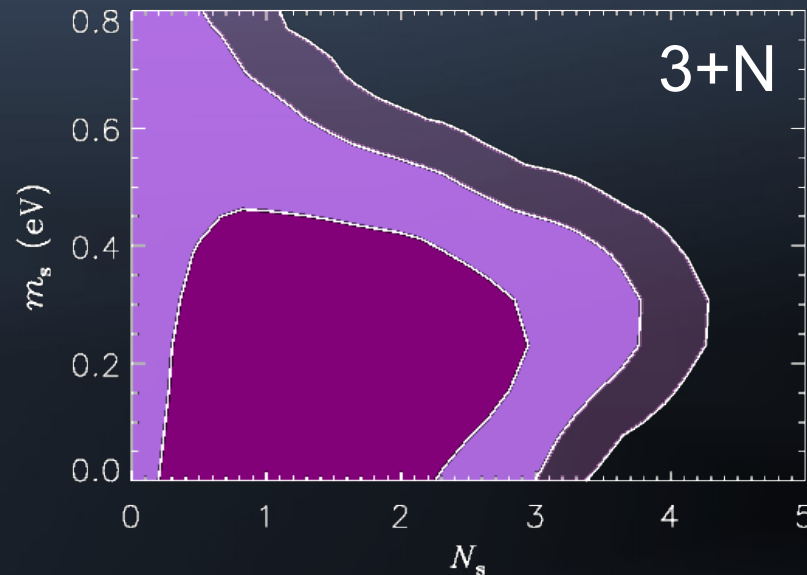
Two qualitatively different mass hierarchies:



[JH, Hannestad, Raffelt, Tamborra, Wong (2010)]

Sterile neutrino scenario

- ♦ 3+1, 3+2, 3+3 are fine ...
... *as long as the steriles are light enough!*
- ♦ Unfortunately, 1 eV appears to be somewhat too heavy
- ♦ Reminder: we assumed minimal extension of cosmological standard model (Λ CDM + N_{eff} + m_ν)



[JH, Hannestad, Raffelt, Tamborra, Wong (2010)]

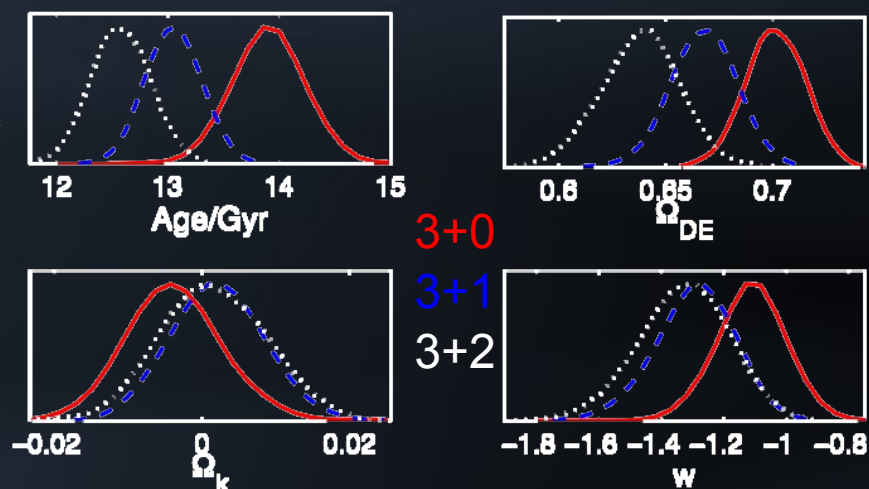
Impact on cosmological model

- Assume laboratory hints for steriles are real, fix masses to best-fit values in 3+1/3+2

| | Δm_{41}^2 [eV ²] | $ U_{e4} $ | Δm_{51}^2 [eV ²] | $ U_{e5} $ | χ^2/dof |
|-----|--------------------------------------|------------|--------------------------------------|------------|---------------------|
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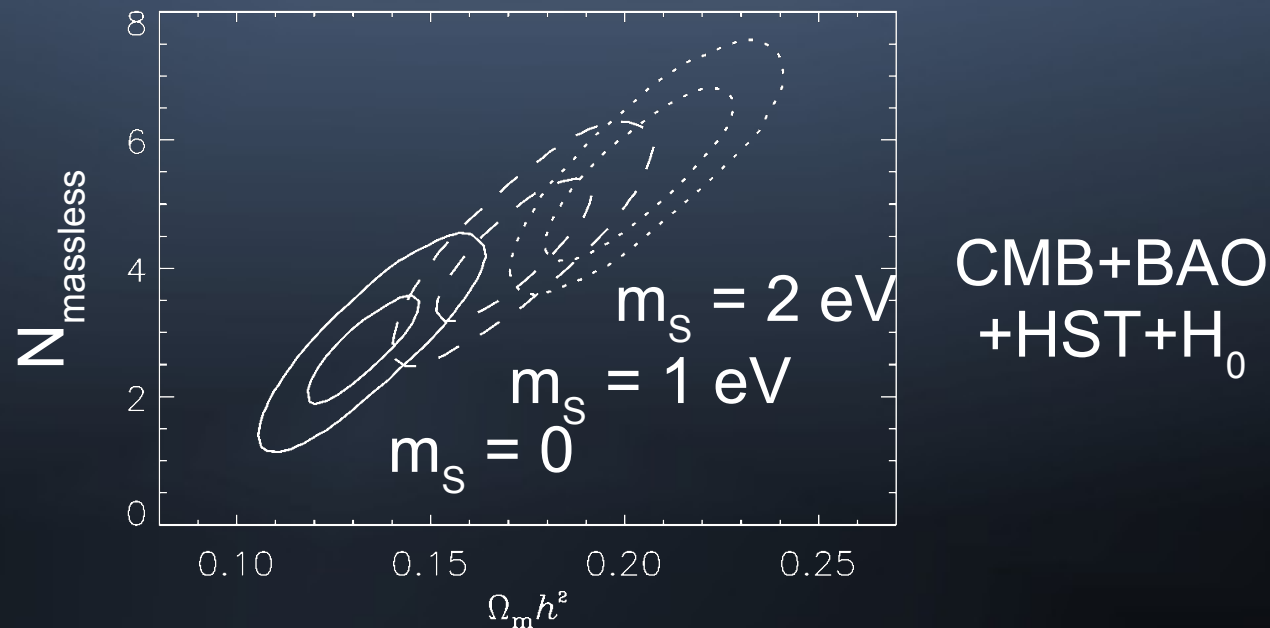
Table I: Best fit points for the 3+1 and 3+2 scenarios from reactor anti-neutrino data.

- Extend model and allow curvature parameter Ω_k and dark energy equation of state parameter w to vary
- $w < -1$ at more than 95% c.l.



[Kristiansen, Elgarøy (2011)]

Impact on cosmological model



- ◆ E.g., 1 massive + N massless species
- ◆ For eV-mass steriles: prefer additional massless species and high matter density

[JH, Hannestad, Raffelt, Wong (in preparation)]



PLANCK

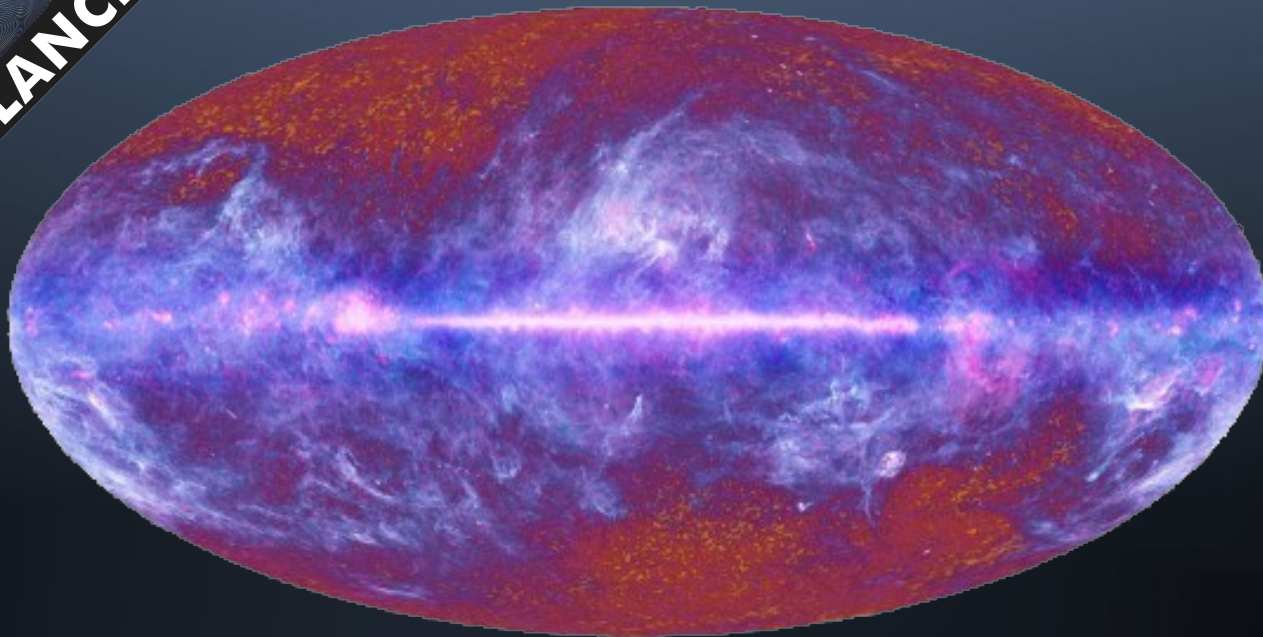


- ◆ Launched 9th May 2009
- ◆ In orbit around Lagrange point L₂
- ◆ Measures CMB in 9 frequency channels 30-857 GHz
- ◆ ~ 5 arcmin resolution
 - limited by cosmic variance up to multipoles of ~2000
- ◆ Expected sensitivity to N_{eff} : $\sigma_{N_{\text{eff}}} \approx 0.2$

[JH, Lesgourgues, Mangano (2007)]



PLANCK



- ♦ January: PLANCK early papers (mostly concerning instrument performance and foreground physics)
- ♦ Cosmology papers: early 2013

Conclusions

- ◆ Cosmological data show slight preference for additional relativistic degrees of freedom, such as sterile neutrinos
- ◆ 3+2 sterile neutrino interpretation of LSDN/MiniBooNE/reactor data is problematic if implemented in naïve minimal cosmological model (too many for BBN, too heavy for CMB+LSS)
- ◆ Incomplete thermalisation or an extension of the cosmological model are required for compatibility
- ◆ Exciting results from PLANCK soon!

